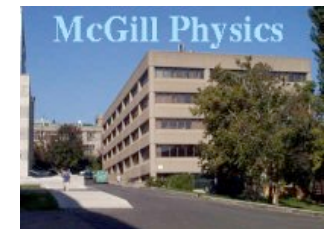
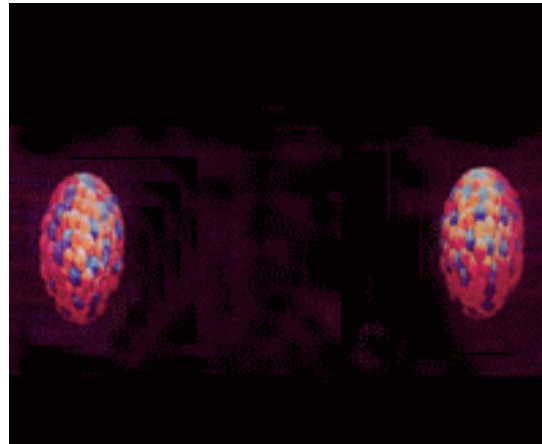


International Conference on Topics in Heavy Ion Collisions

Montreal, 25-28 June 2003

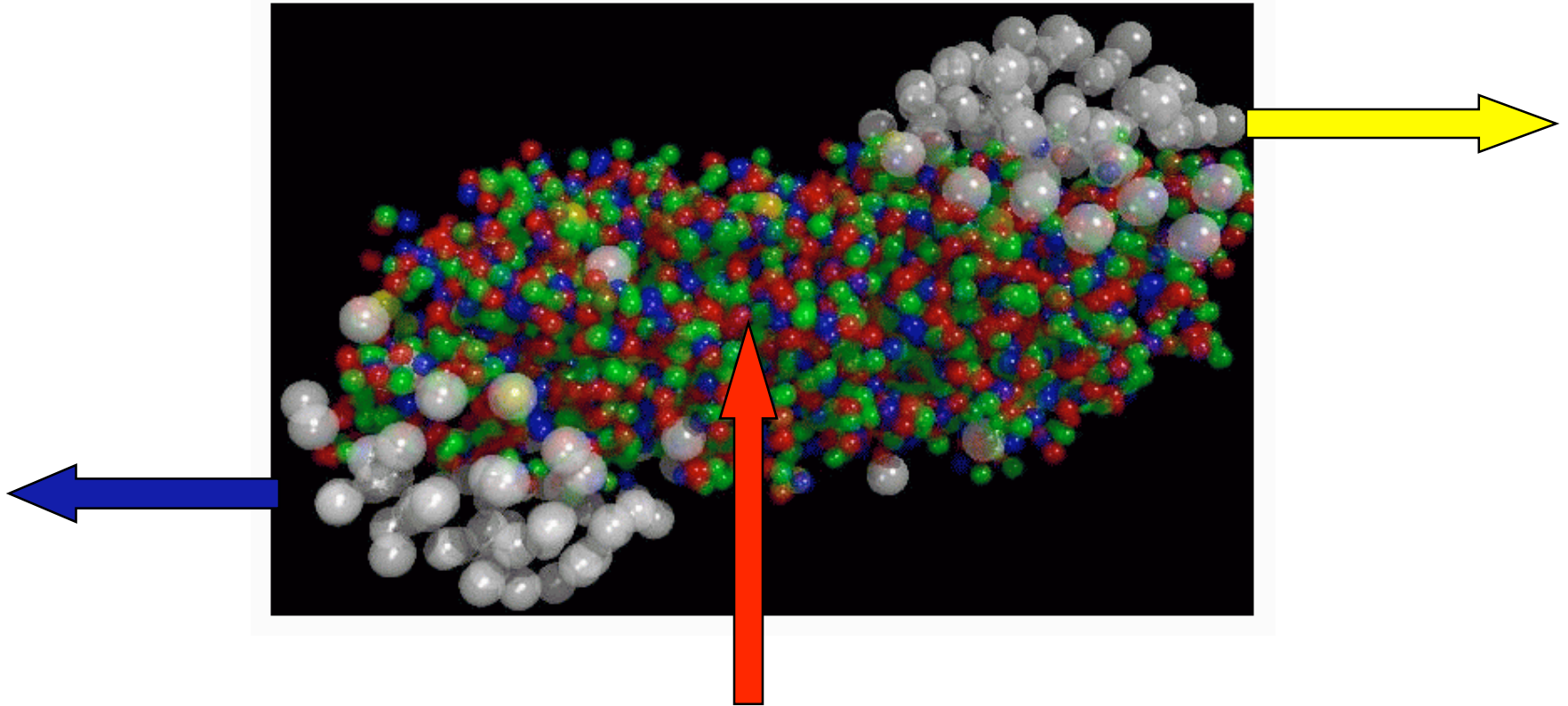
Chiral Condensate Dynamics

Jørgen Randrup, LBNL, Berkeley



Topics in Heavy Ion Collisions, Montreal 25-28 June 2003

[UrQMD: 60 GeV/N Au + 60 GeV/N Au]



Central rapidity region:

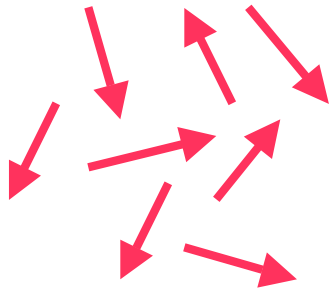
- * Highly **excited** strongly interacting matter
- * Rapid **expansion** (primarily longitudinally)

QCD \Rightarrow Chiral Condensate

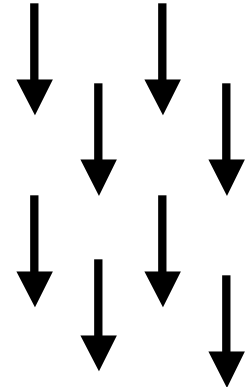
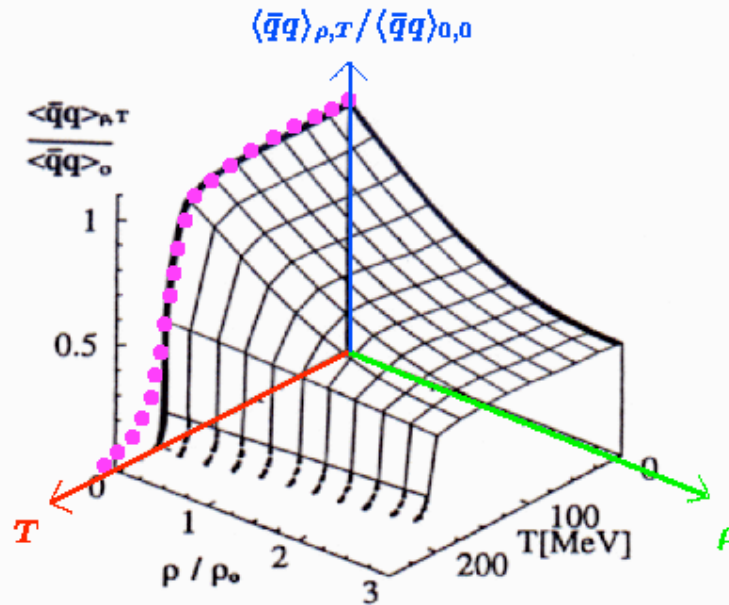
$$q = \begin{pmatrix} u \\ d \end{pmatrix} : \mathcal{L}_{QCD} = \bar{q}(i\gamma^\mu D_\mu - m)q - \frac{1}{4}F_{\mu\nu}^a F_a^{\mu\nu}$$

The QCD vacuum is a condensate of quark-antiquark pairs:

$$\langle \bar{q}q \rangle \equiv \langle u^\dagger \gamma^0 u + d^\dagger \gamma^0 d \rangle \neq 0$$



Hot



Cold

[T. Waas, R. Brockmann and W. Weise: Phys. Lett. B405 (1997) 215]

Chiral Condensate Dynamics



Supercritical fields?

$$\mu_{\pi}^2 < 0$$



Disoriented chiral condensates? $\langle \bar{q}\tau\gamma_5 q \rangle \neq 0$



Chiral antenna?

$$\mu_{\pi}(t)$$



“Emergence of coherent long-wavelength oscillations after a quench”

[K. Rajagopal & F. Wilczek: Nucl. Phys. B404 (1993) 577]

$$\Phi(\mathbf{r}, t) = (\phi, \vec{\pi}): \quad L = \int d^3r dt \left[\frac{1}{2} \partial^i \phi \partial_i \phi + \frac{1}{4} \phi (\partial^i \vec{\pi} \partial_i \vec{\pi} + v^2)^2 + H \phi \right]$$

Eq of motion
for pion field

$$\frac{d^2}{dt^2} \phi_k(t) = -\omega_k^2(t) \phi_k(t) = -[k^2 + \omega_\pi^2(t)] \phi_k(t)$$

Effective
pion mass

$$\omega_\pi^2(t) = \phi[\omega_\pi^2(t) + v^2]$$



When

$$\omega_\pi^2 < v^2$$

then

$$\omega_\pi^2 < 0$$

so

$$\phi_k \sim e^{\pm \int_0^t dt' |\omega_k(t')|}$$

Quench

=>

Super-critical

=>

Explosive growth!

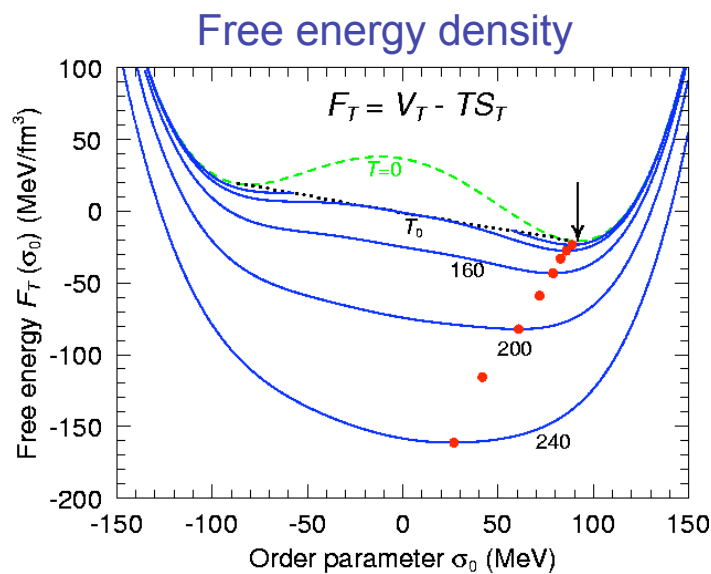
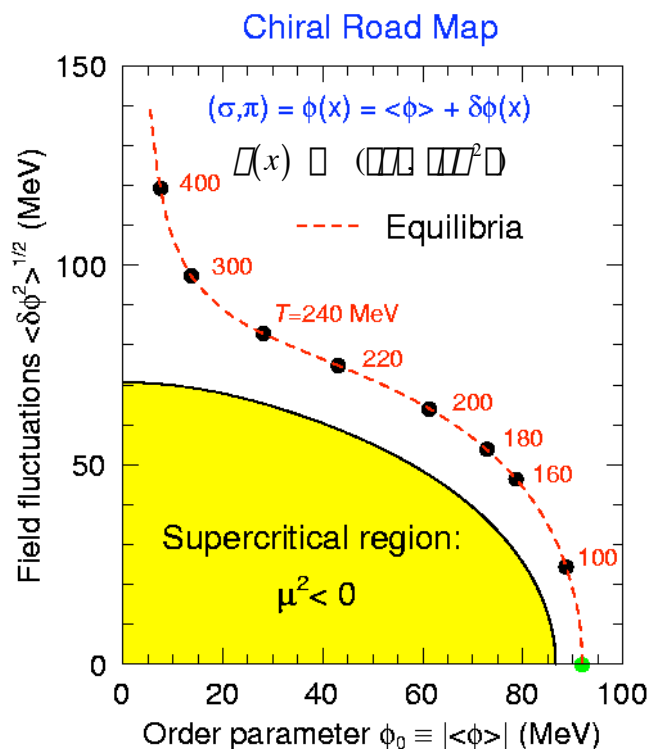
Spontaneous
pion-pair
production

How to analyze a dynamical chiral field?

EoM:
$$\left[\partial_t^2 \sigma + \sigma(\sigma^2 + v^2) \right] \sigma_0 = He_\sigma \quad \Rightarrow \quad \sigma(\mathbf{r}, t) = \left(\sigma(\mathbf{r}, t), \vec{\pi}(\mathbf{r}, t) \right)$$

Project the dynamical state onto a chiral phase diagram

$$\left\{ \begin{aligned} \langle \sigma \rangle &= \int d^3r \sigma(\mathbf{r}, t) \\ \langle \pi^2 \rangle &= \int d^3r (\sigma(\mathbf{r}, t) - \langle \sigma \rangle)^2 \end{aligned} \right.$$



$$m_\sigma^2 = \left[\frac{\partial^2 \mathcal{L}}{\partial \sigma^2} + \frac{\partial^2 \mathcal{L}}{\partial \pi^2} + 2 \frac{\partial^2 \mathcal{L}}{\partial \sigma \partial \pi} v^2 \right]$$

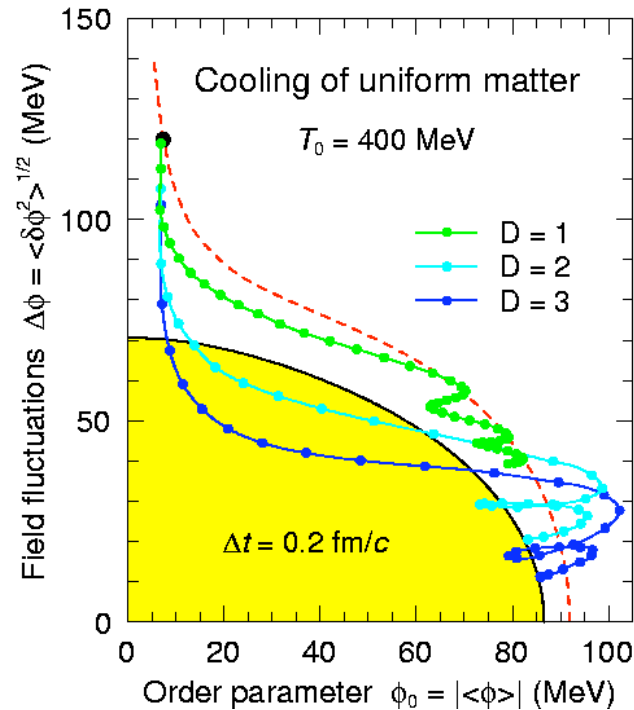
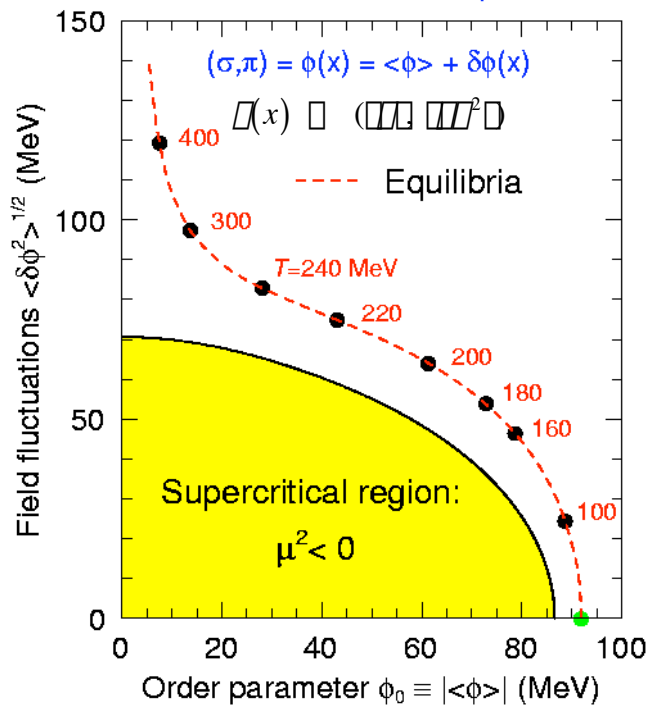
Effective mass m_σ

Will a quench occur?

Emulate a D-dimensional Bjorken scaling expansion of a hot uniform system

$$\left[\partial_t^2 \phi + \lambda (\phi - \phi_0 - v^2) \right] \phi - H e_\phi = \lambda \frac{D}{t} \partial_t \phi$$

Chiral Road Map



Supercriticality is NOT likely to occur!

Chiral Condensate Dynamics

Supercritical fields?

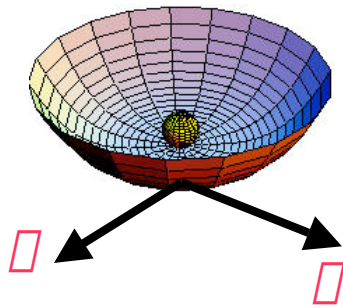


The early expansion is predominantly longitudinal $\Rightarrow \mu_{\pi} > 0$

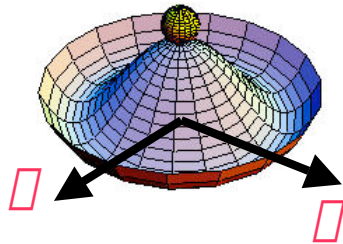
Disoriented chiral condensates?



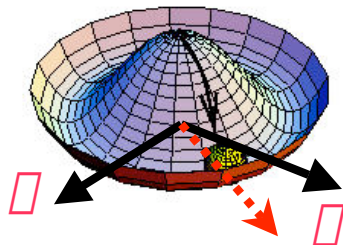
Disoriented Chiral Condensates?



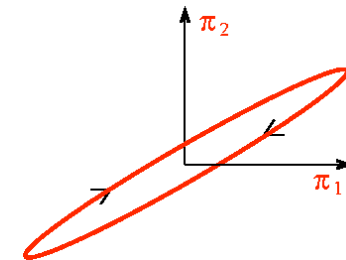
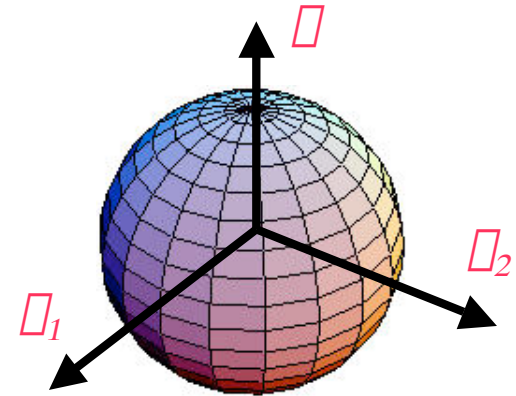
The environment is hot:
 $O(4)$ symmetry is favored



The environment cools:
 $O(4)$ symmetry disfavored



The order parameter grows
in a disoriented direction

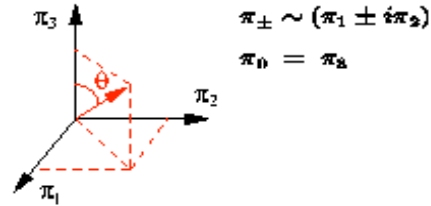


Isospin-directed
field oscillations

Neutral pion fraction: very idealized

Pion field: $\vec{\pi} = (\pi_1, \pi_2, \pi_3)$

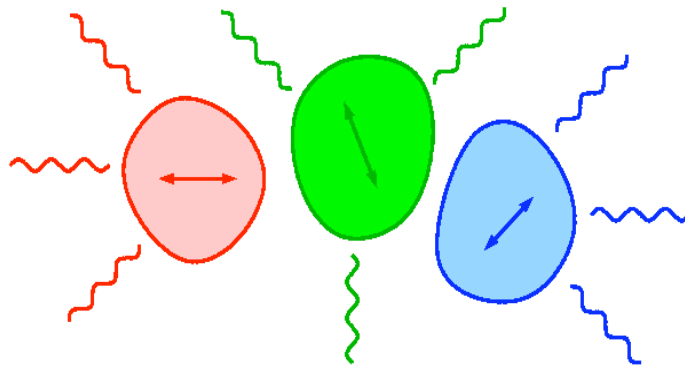
$$n_0 \sim |\pi_3|^2 = |\vec{\pi}|^2 \cos^2 \theta$$



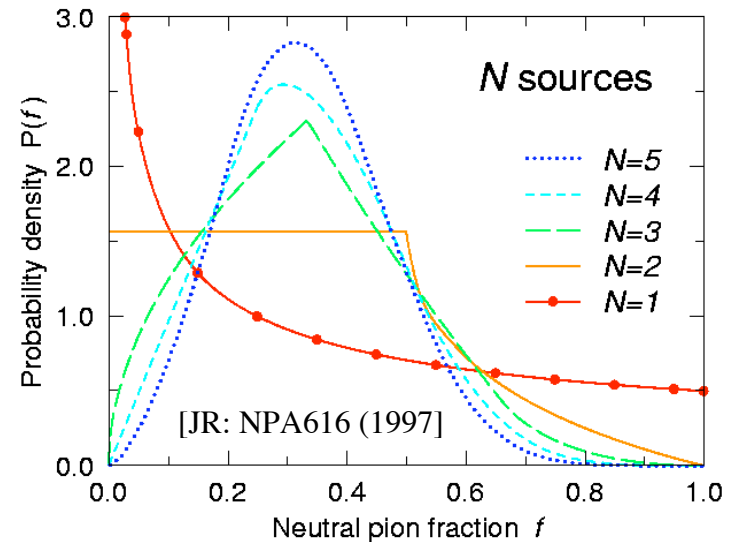
$$\text{Neutral pion fraction : } f = \frac{n_0}{n_- + n_0 + n_+} = \cos^2 \theta$$

[AA Amsden & MG Ryskin:
Phys Lett B266 (1991) 482]

$$P(f)df = d \cos \theta \Rightarrow P(f) = \frac{d \cos \theta}{df} = \left(\frac{df}{d \cos \theta} \right)^{-1} = (2 \cos \theta)^{-1} = \frac{1}{2\sqrt{f}}$$

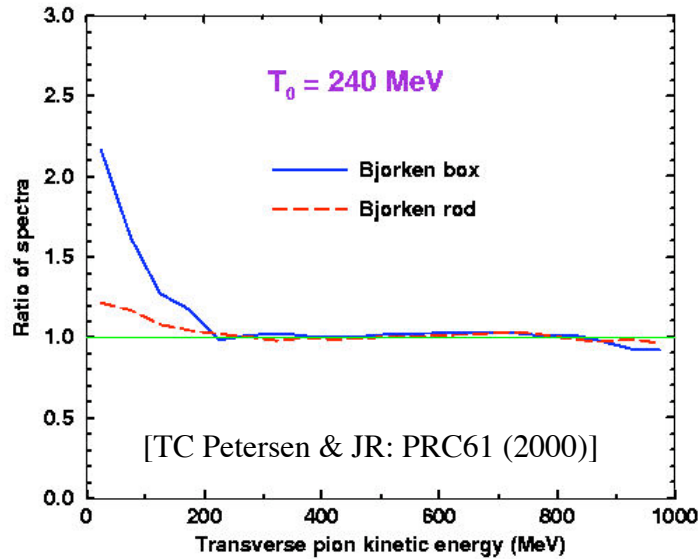


Domains of Disoriented Chiral Condensates

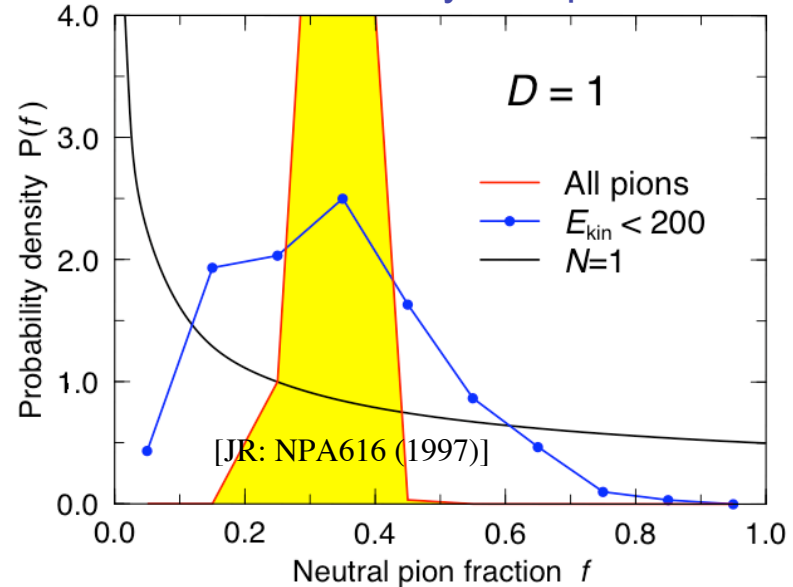


Neutral pion fraction: less idealized

Transverse pion spectrum:



From all to only soft pions:



=>

The chiral dynamics affects only the soft pion modes!

$\pi^0 \rightarrow 2\pi$ cannot be cut experimentally!

Experimental searches have found no anomaly in $P(f_{\pi^0})$:



MiniMax Collaboration: PRD55 (1997)



WA98 Collaboration: PLB420 (1998)

Chiral Condensate Dynamics



Supercritical fields?

The early expansion is predominantly longitudinal $\Rightarrow \mu_\pi > 0$



Disoriented chiral condensates?

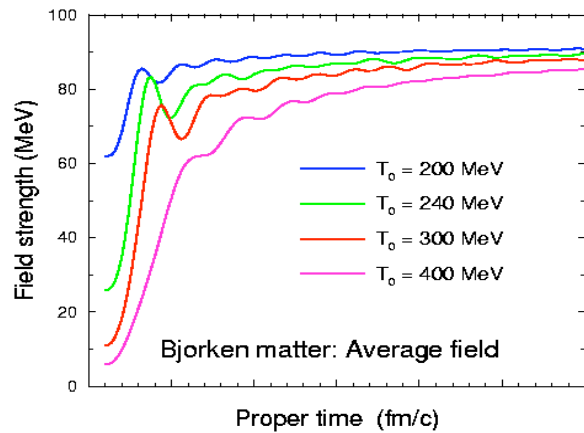
Measuring the *soft* neutral pion fraction is impractical: $\pi^0 \rightarrow 2\gamma$



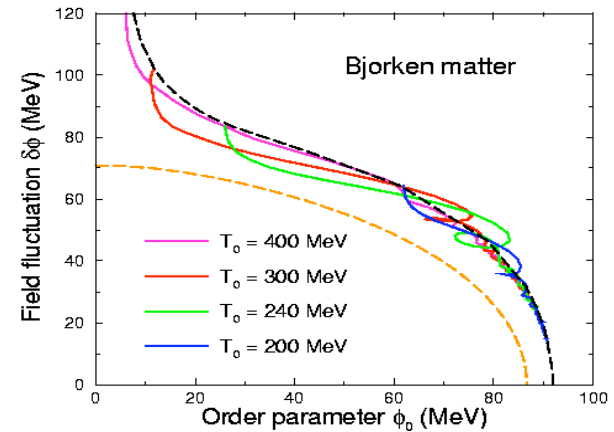
Chiral antenna?

Time-dependent chiral environment

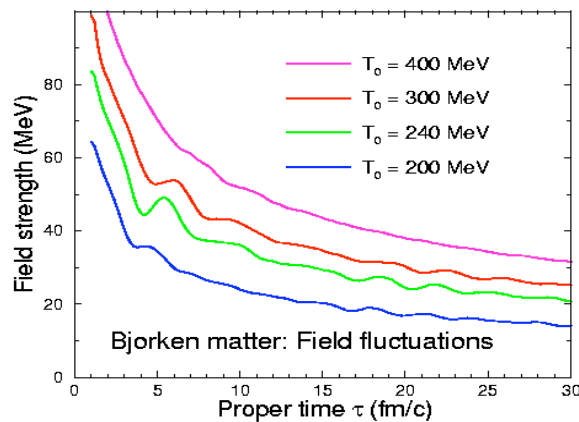
Order parameter $\langle \sigma \rangle$



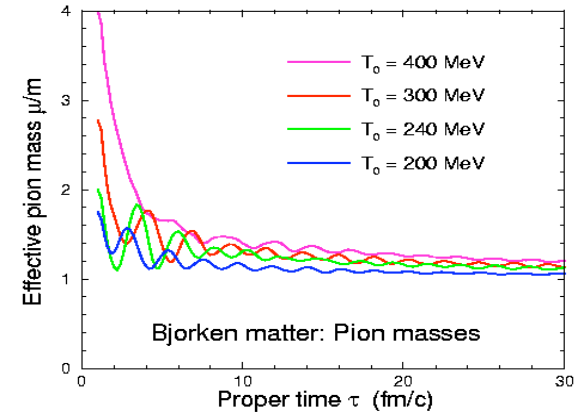
Phase evolution (σ, π)



Fluctuation $\langle \sigma^2 \rangle$



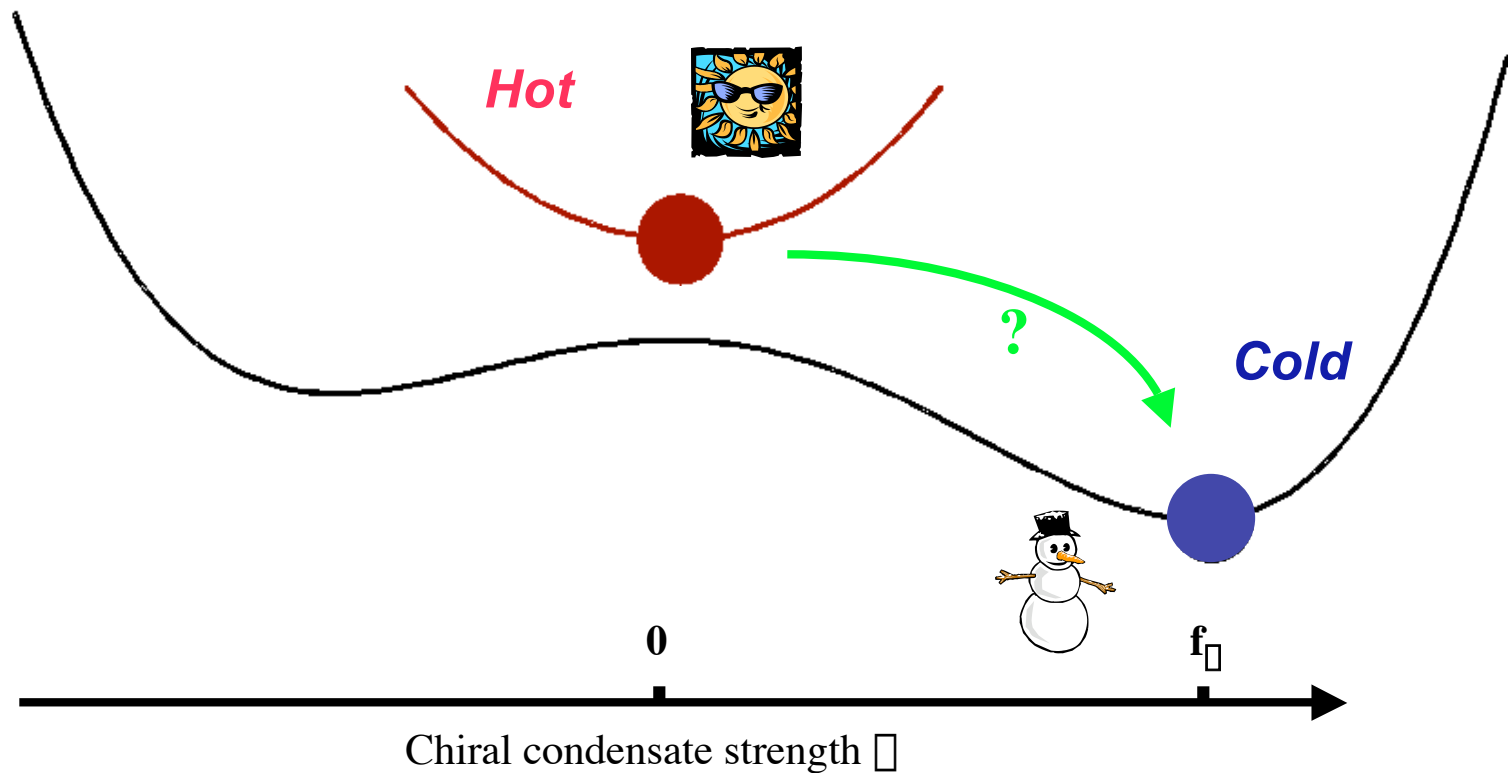
Effective mass $m_\pi(t)$

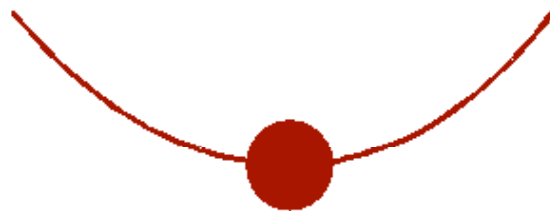


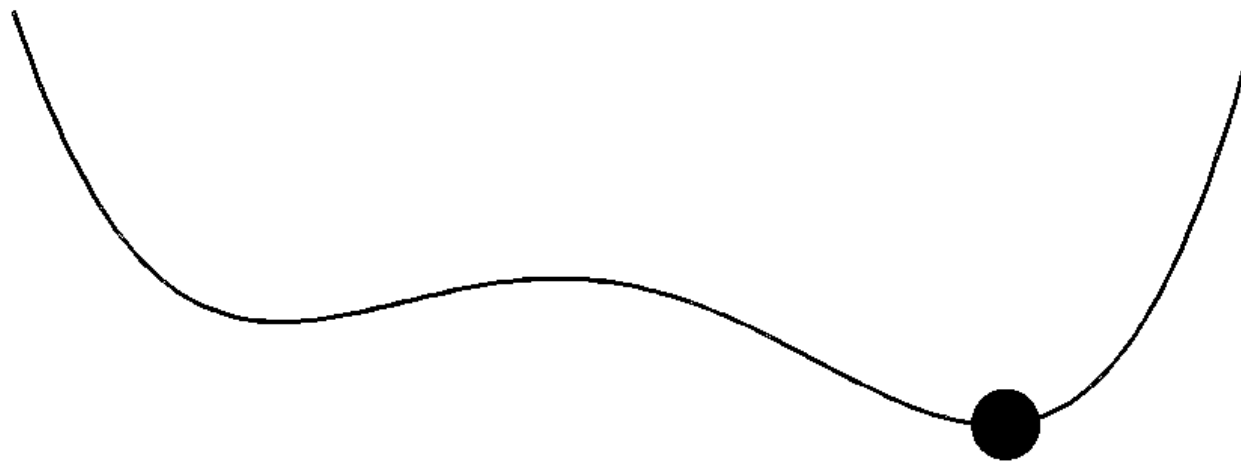
$$m_i^2(t) = m^2 \left[\sigma_0^2 + \langle \pi^2 \rangle + 2 \langle \pi \pi_i \rangle v^2 \right]$$

Chiral condensate dynamics: Non-equilibrium relaxation

The effective potential evolves from hot to cold form, bringing the condensate out of equilibrium:





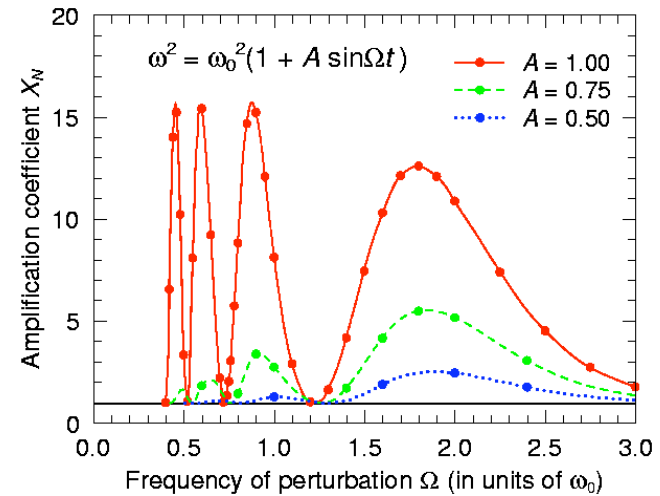


Frequency modulation => pair creation

$$\omega^2(t) \equiv \omega^2[k^2 + v^2] : \quad \omega^2(t) = k^2 + v^2$$



$$\left[\partial_t^2 + \omega^2(t) \right] \varphi(t) = 0$$



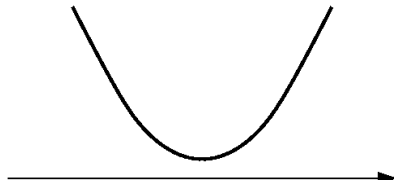
Mean particle number:
vacuum fluctuations help

$$\langle n \rangle \equiv \langle a^\dagger a \rangle : \quad \langle n(t) \rangle = X(t) \left[n(0) + \frac{1}{2} \right] \approx \frac{1}{2}$$

Number fluctuations
are enhanced

$$\langle n^2 \rangle \equiv \langle n^2 \rangle - \langle n \rangle^2 : \quad \begin{cases} n(0) = 0 : \langle n^2(t) \rangle = 2 \langle n(t) \rangle \langle \bar{n}(t) \rangle \\ n(0) \ll 1 : \langle n^2(t) \rangle \approx X(t)^2 \left[\langle n^2(0) \rangle + \frac{1}{2} \right] \approx \frac{1}{2} \end{cases}$$

The number of quanta changes by **two** at a time!



$$H(t) = \frac{1}{2}p^2 + \frac{1}{2}\omega^2(t)q^2$$

$$\begin{cases} q = \frac{1}{\sqrt{2\omega_0}}[a + a^\dagger] \\ p = i\sqrt{\frac{\omega_0}{2}}[a - a^\dagger] \end{cases}$$

$$H(t) = \frac{1}{2}\omega_+(t)[aa^\dagger + a^\dagger a] + \frac{1}{2}\omega_-(t)[a^2 + (a^\dagger)^2]$$

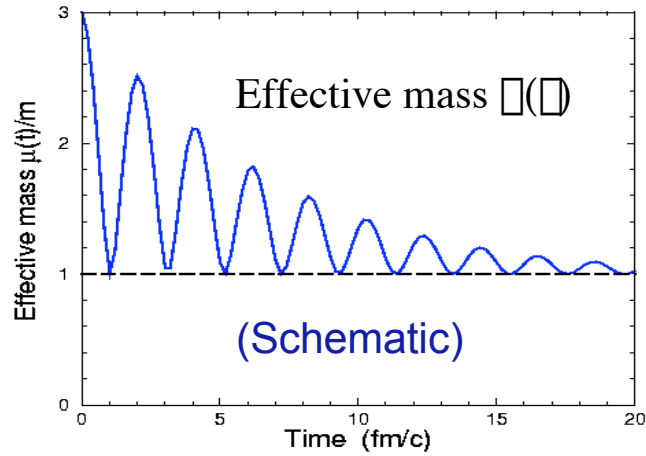
$$\omega_\pm(t) = [\omega^2(t) \pm \omega_0^2]/2\omega_0^2$$

OBS: The particles are created pairwise

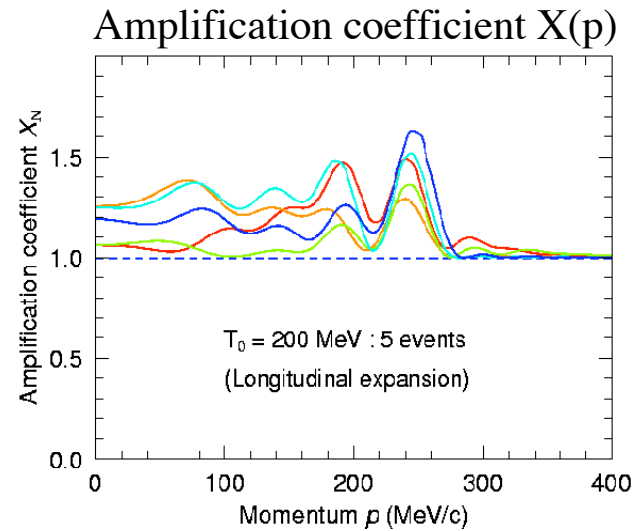
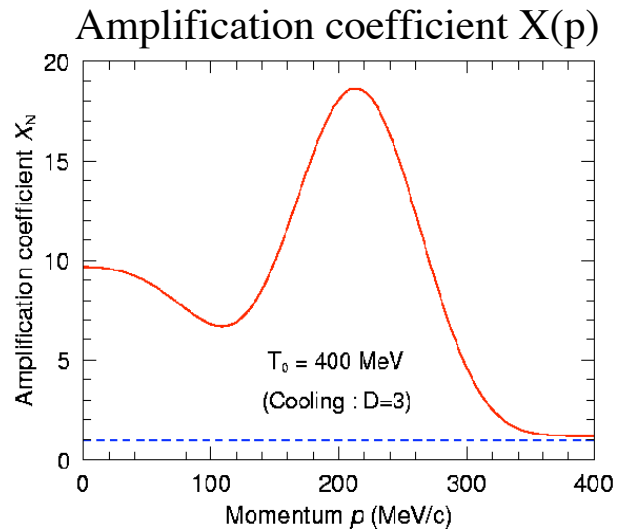
In a wave mechanical picture, this is a direct consequence of the reflection symmetry, which conserves the wave function's parity:

An even wave function will remain even and an odd wave function will remain odd

General character of the mass evolution



- ★ Expansion: overall decrease of μ
- ★ Non-equilibrium: oscillations of μ



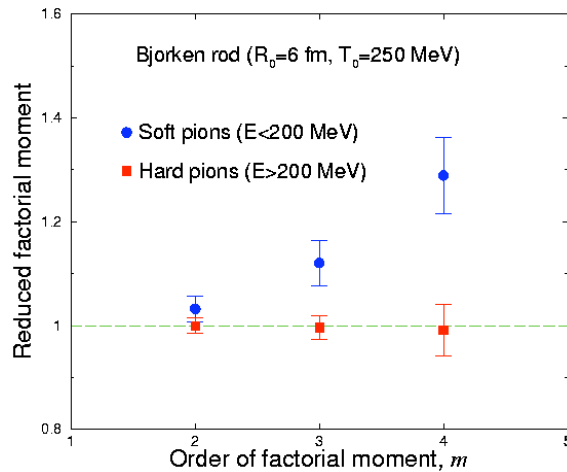
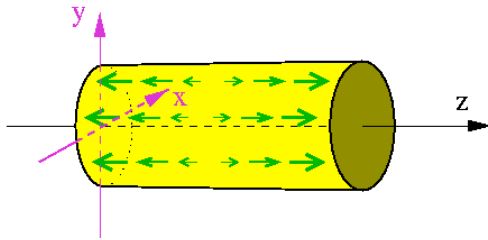
Fluctuations in the number of soft pions

Factorial moments of the multiplicity distribution:

$$M_m \equiv \langle N(N-1)\cdots(N-m+1) \rangle$$

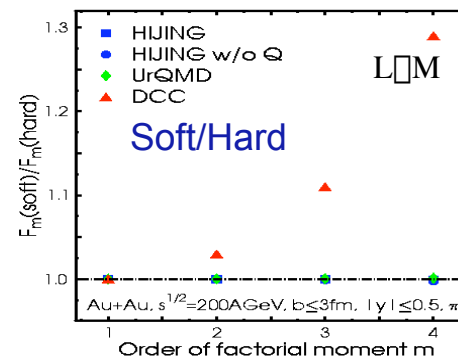
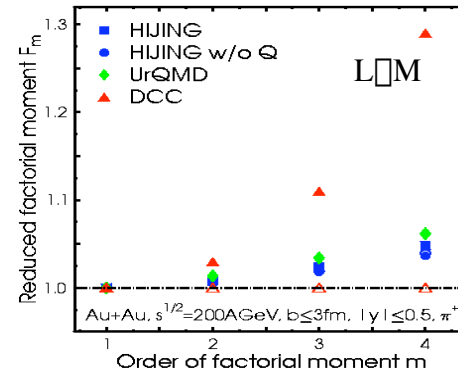
-- are enhanced by the time modulation

Linear \square model simulations for a stretching cylinder:



[TC Petersen & JR: PRC61 (2000)]

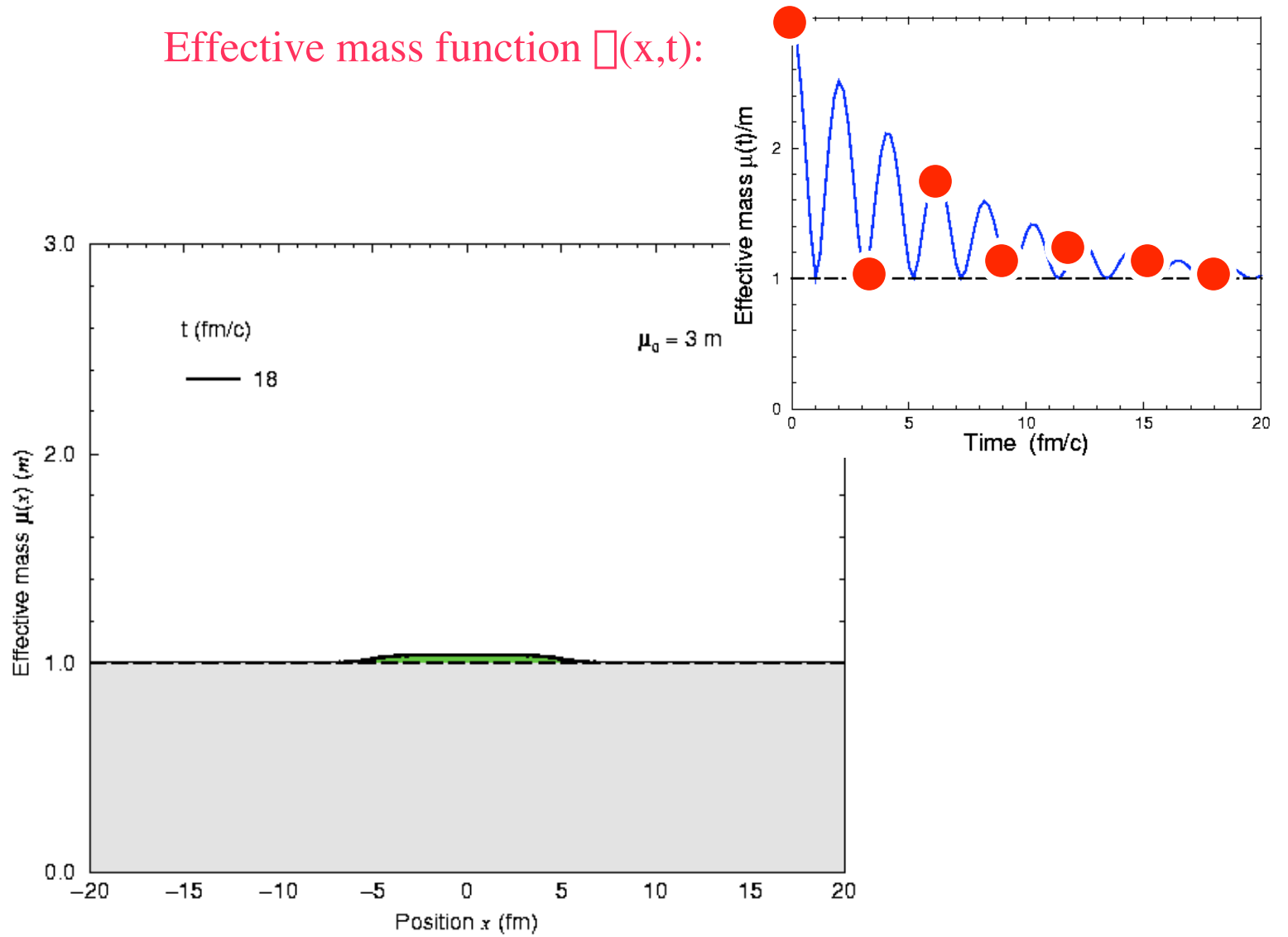
Comparison with HIJING & UrQMD:



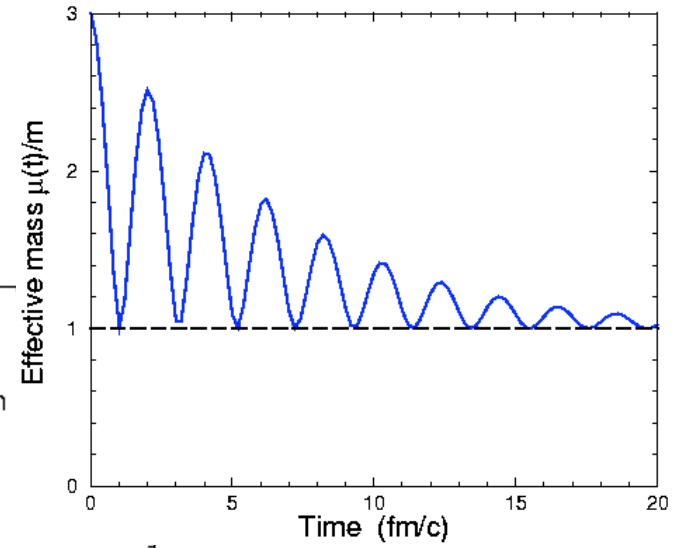
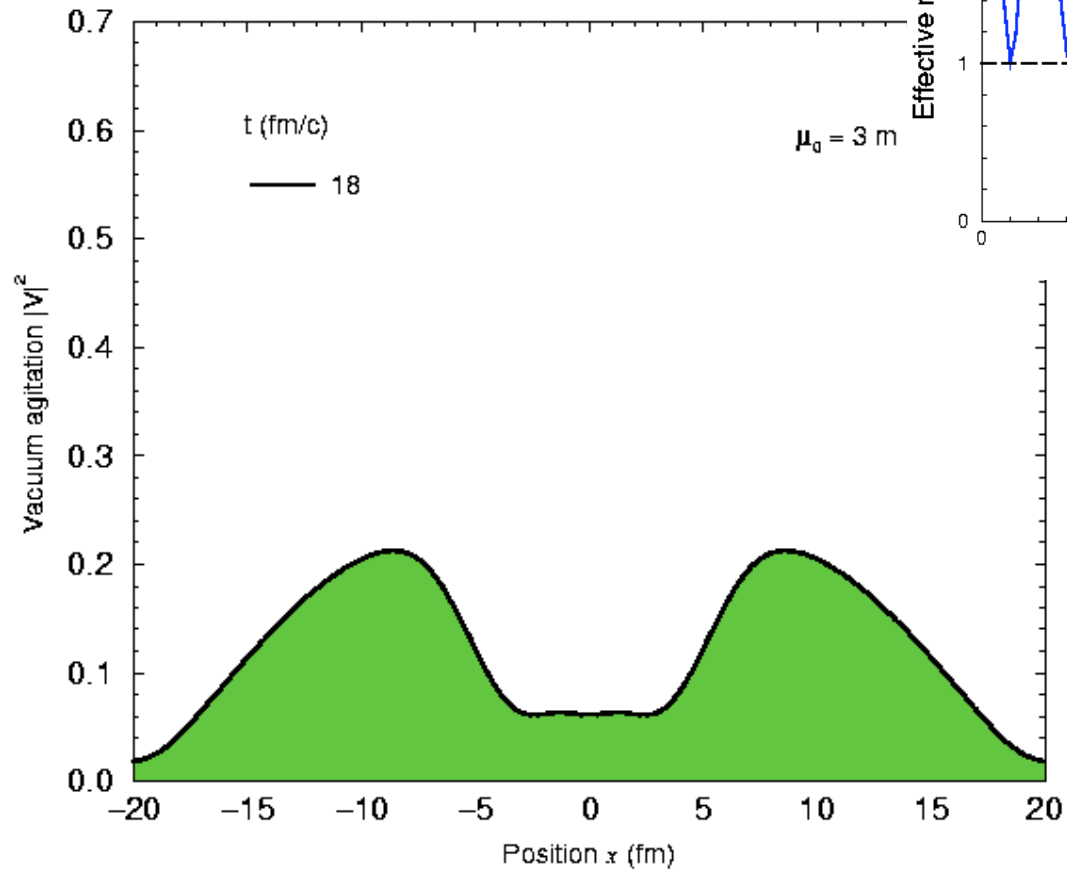
<-- Soft = Hard

[M Bleicher et al: PRC62 (2000)]

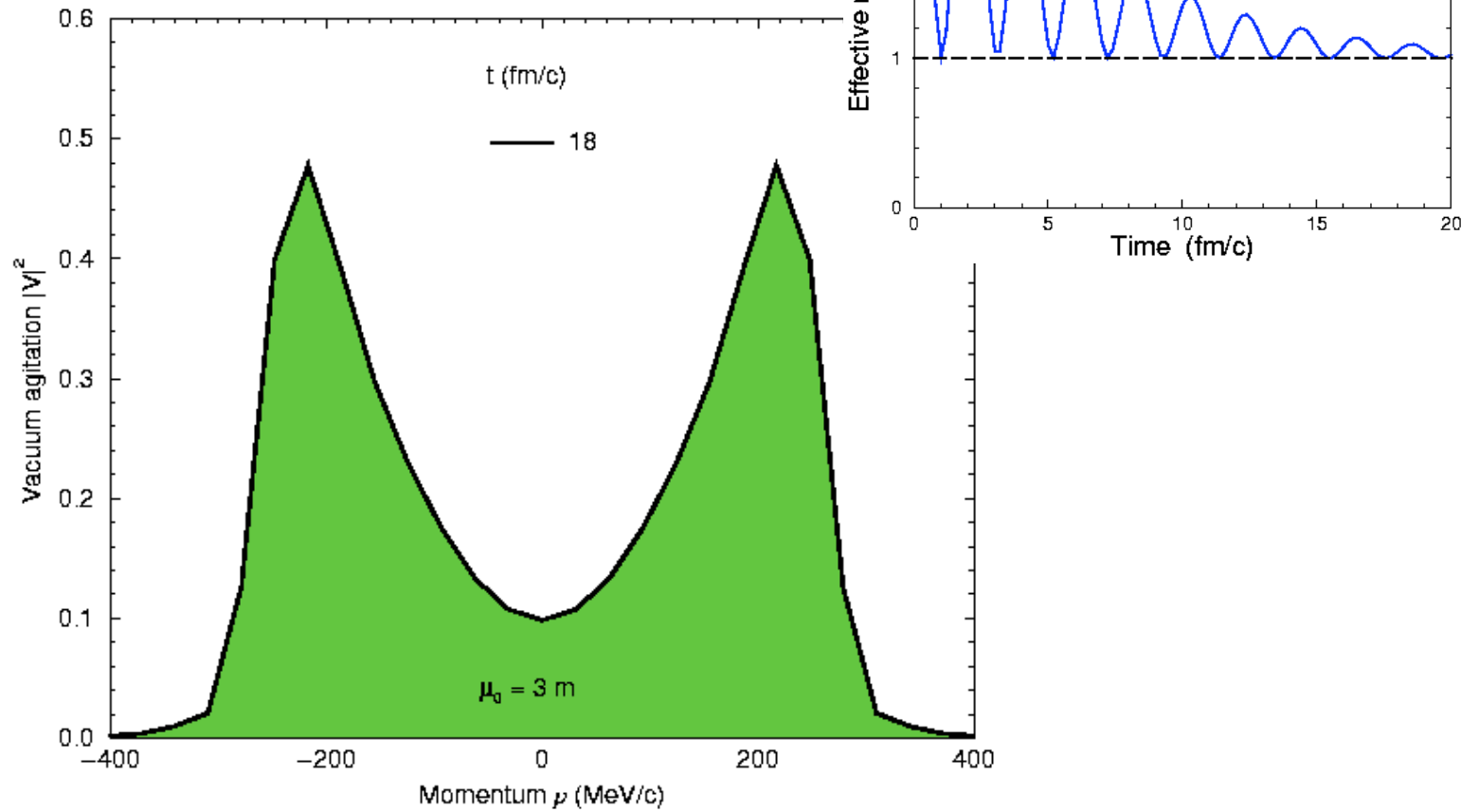
Effective mass function $\mu(x,t)$:



Spatial density of quasiparticles:

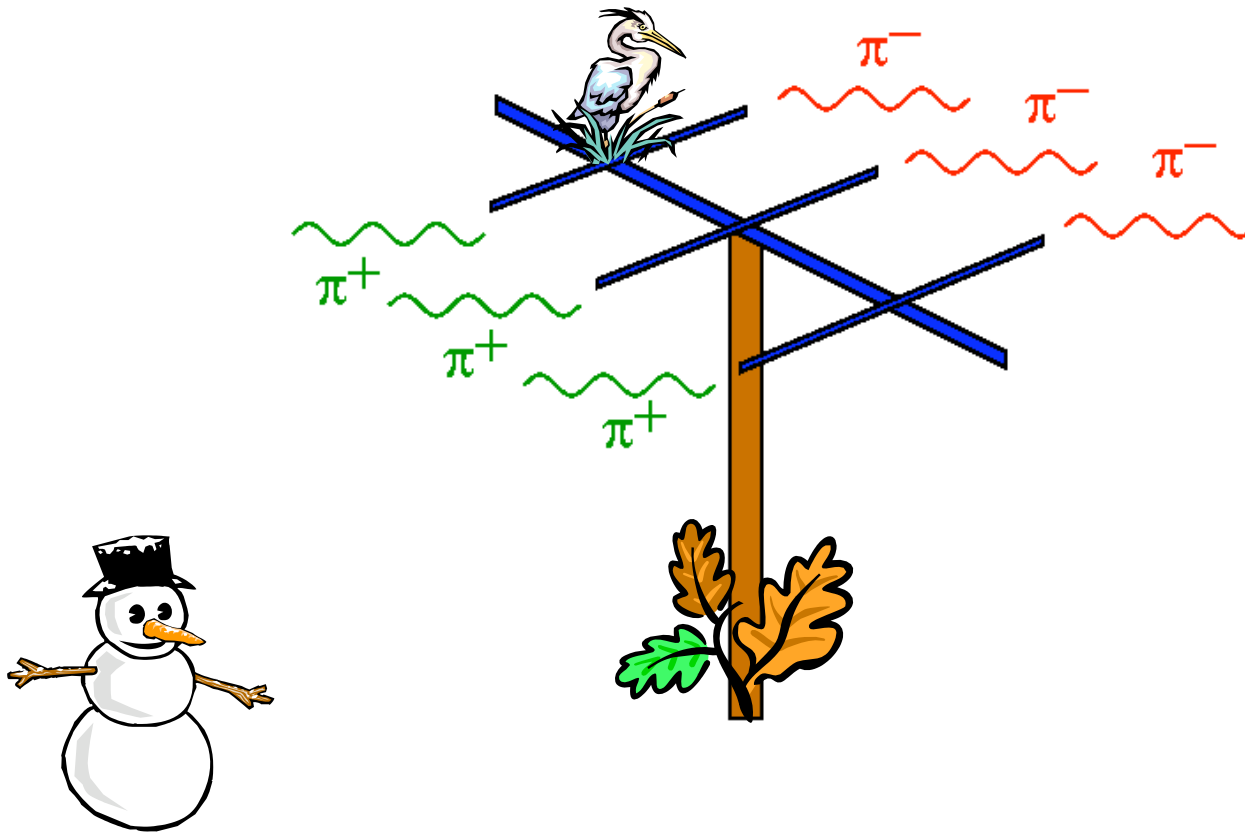


Momentum spectrum of quasiparticles:



Chiral antenna

Soft charge-conjugate back-to-back pion pairs



Relaxation dynamics of the chiral order parameter

[J. Randrup, *Physical Review C* 63 (2001) 061901(R); 65 (2002) 054906]

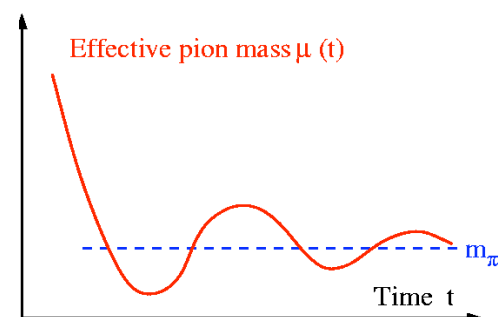
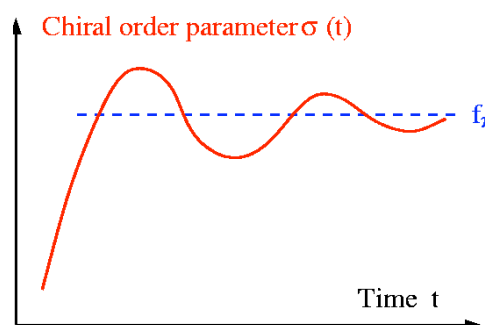
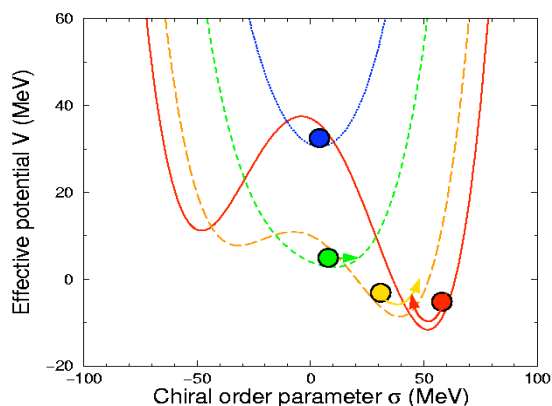
The potential changes fast from hot to cold

=>

The order parameter shows non-equilibrium relaxation

=>

The pion mass exhibits an oscillatory evolution

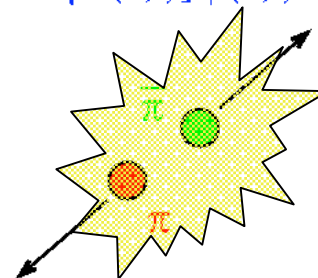


The rapid change of V leads to non-equilibrium relaxation of σ

The pionic degrees of freedom see an undulating environment

Emission of **soft back-to-back charge-conjugate pion pairs**

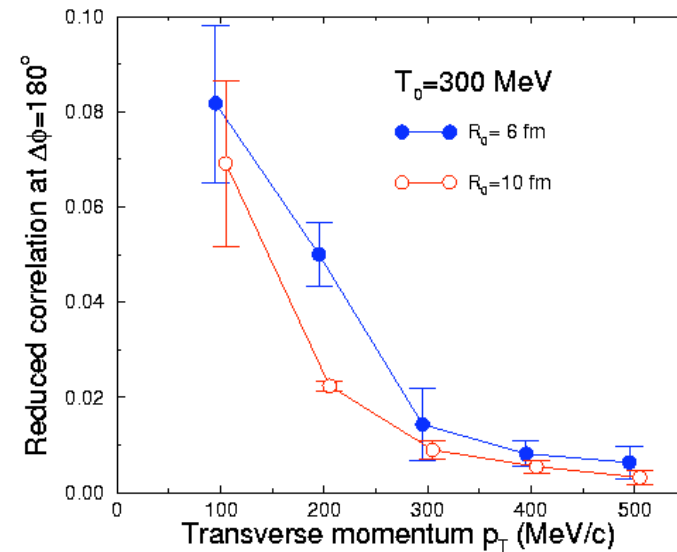
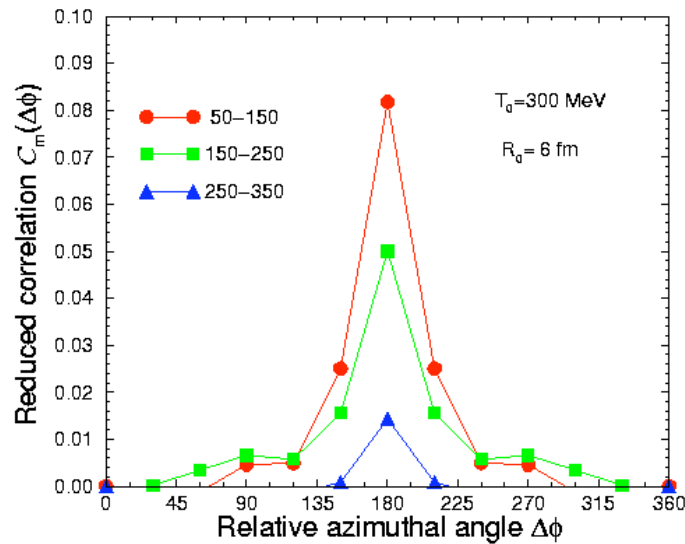
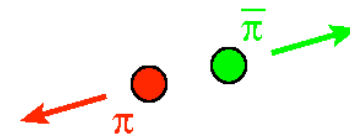
$$[\square + \mu^2(x,t)]\phi(x,t) = 0$$



Correlations among the soft pions

The dynamical change of the effective mass creates back-to-back pairs of soft charge-conjugate pions:

$$[\square + \mu^2(x,t)]\phi(x,t) = 0$$



Chiral Antenna



The antenna signal is **sensitive** to the properties of interest, such as the initial degree of chiral symmetry restoration and the subsequent rate of expansion/cooling.
So, if identified, it may be a useful probe of these features.

But the calculations should **not** be taken at face value,
as they are subject to a number of caveats:



Signal strength: additional background?

LQM underpredicts the pion multiplicity => Correlation signal is reduced



Competing mechanisms: Meson decays?

$K_s^0(500) \rightarrow 2\pi$	$\rho(770) \rightarrow 2\pi$	$\omega(780) \rightarrow 3\pi$
just above threshold	well above threshold	not back-to-back



Signal degradation: collisions?

Baked Alaska (Bjorken): Signal pions are **late** & **slow**, so don't collide
[Input from transport models needed!]



Linear \square model is inadequate?

Chiral condensate dynamics

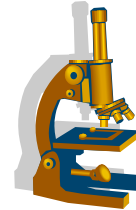


Novel phenomena:

Supercritical strong field?

Isospin-directed chiral condensates?

Parametric amplification?



Observational consequences:

Spontaneous pion pair creation

Anomalous distributions of neutral π (and K) fractions

Enhanced fluctuations and correlations



Status:

