

Thermal Conformal Blocks

Yan Gobeil

McGill University

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Thermal Conformal Blocks

Yan Gobeil,^{a,*} Alexander Maloney,^{a,○} Gim Seng Ng,^{b,c,◇} Jie-qiang Wu^{d,▷}

^a*Department of Physics, McGill University, Montréal, QC, Canada*

^b*School of Mathematics, Trinity College Dublin, Dublin 2, Dublin, Ireland*

^c*Hamilton Mathematical Institute, Trinity College Dublin, Dublin 2, Ireland*

^d*Center for Theoretical Physics, Massachusetts Institute of Technology, Cambridge, MA 02139, USA*

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Diagram in
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CFT = QFT with scale invariance

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CFT = QFT with scale invariance

Symmetry generators:

- Translations P_μ , Rotations $M_{\mu\nu}$
- Dilatations D , Special Conformal Transformations K_μ

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Irreps of conformal group = Conformal families (HW reps)

- Primary $|\mathcal{O}\rangle$: $K_\mu |\mathcal{O}\rangle = 0$, $D |\mathcal{O}\rangle = \Delta_{\mathcal{O}} |\mathcal{O}\rangle$, spin
- Descendants $|\mathcal{O}, n\rangle \sim P^n |\mathcal{O}\rangle$, $D |\mathcal{O}, n\rangle = (\Delta_{\mathcal{O}} + n) |\mathcal{O}, n\rangle$

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Study scalar field $\phi(x)$ in thermal state with $T = 1/\beta$

$$\begin{aligned}\langle \phi(x) \rangle_\beta &= \text{Tr}[\phi(x)e^{-\beta H}] = \text{Tr}[\phi(x)e^{-\beta D}] = \sum_i q^{\Delta_i} \langle i | \phi(x) | i \rangle \\ &= \sum_{\text{primaries } \mathcal{O}} q^{\Delta_{\mathcal{O}}} \sum_n \langle \mathcal{O}, n | \phi(x) | \mathcal{O}, n \rangle q^n\end{aligned}$$

where $q \equiv e^{-\beta}$ and assume orthonormal basis for simplicity.

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where $q \equiv e^{-\beta}$ and assume orthonormal basis for simplicity.

Conformal symmetry relates $\langle \mathcal{O}, n | \phi | \mathcal{O}, n \rangle$ to $\langle \mathcal{O} | \phi | \mathcal{O} \rangle$ so define

$$\langle \phi \rangle_\beta = \sum_{\text{primaries } \mathcal{O}} \langle \mathcal{O} | \phi | \mathcal{O} \rangle q^{\Delta_{\mathcal{O}}} F(\Delta_\phi, \Delta_{\mathcal{O}}, \ell_{\mathcal{O}}; q)$$

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Purely kinematic object, focus on $\ell_{\mathcal{O}}$

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- Conformal algebra:

$$[D, P_\mu] = P_\mu, \quad [D, K_\mu] = -K_\mu, \quad [K_\mu, P_\nu] = 2\delta_{\mu\nu}D - 2M_{\mu\nu}$$

$$[M_{\mu\nu}, P_\rho] = \delta_{\nu\rho}P_\mu - \delta_{\mu\rho}P_\nu \quad , \quad [M_{\mu\nu}, K_\rho] = \delta_{\nu\rho}K_\mu - \delta_{\mu\rho}K_\nu$$

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- For scalars:

$$[D, \phi(x)] = (\Delta_\phi + x^\mu \partial_\mu) \phi(x), \quad [P_\mu, \phi(x)] = \partial_\mu \phi(x)$$

$$[K_\mu, \phi(x)] = (2x_\mu \Delta_\phi + 2x_\mu x^\nu \partial_\nu - x^2 \partial_\mu) \phi(x)$$

$$[M_{\mu\nu}, \phi(x)] = (x_\nu \partial_\mu - x_\mu \partial_\nu) \phi(x)$$

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Using $\langle \mathcal{O} | \phi(x) | \mathcal{O} \rangle \sim (x^2)^{-\Delta_\phi/2}$, get

$$\langle \mathcal{O} | K_\mu \phi(x) P_\nu | \mathcal{O} \rangle = \left(2\Delta \delta_{\mu\nu} - \Delta_\phi \delta_{\mu\nu} + \Delta_\phi^2 \frac{x_\mu x_\nu}{x^2} \right) \langle \mathcal{O} | \phi(x) | \mathcal{O} \rangle$$

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Non-orthonormal basis so need

$$\langle \mathcal{O} | K_\mu P_\nu | \mathcal{O} \rangle = 2\Delta \delta_{\mu\nu}$$

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Non-orthonormal basis so need

$$\langle \mathcal{O} | K_\mu P_\nu | \mathcal{O} \rangle = 2\Delta_{\mathcal{O}} \delta_{\mu\nu}$$

$$\Rightarrow \sum_{\mu, \nu} \langle K_\mu \phi(x) P_\nu \rangle (\langle K_\mu P_\nu \rangle)^{-1} = d + \frac{\Delta_\phi (\Delta_\phi - d)}{2\Delta_{\mathcal{O}}}$$

Result

For example in $3d$ we find

$$F(\Delta_\phi, \Delta_\mathcal{O}; q) = a_0 + a_1 q + a_2 q^2 + \dots$$

with

$$a_0 = 1$$

$$a_1 = 3 + \frac{\Delta_\phi(\Delta_\phi - 3)}{2\Delta_\mathcal{O}}$$

$$a_2 = 6 + \frac{\Delta_\phi(\Delta_\phi - 3)[\Delta_\mathcal{O}\Delta_\phi(\Delta_\phi - 3) + 2(8\Delta_\mathcal{O}^2 + \Delta_\mathcal{O} - 2)]}{4(\Delta_\mathcal{O} + 1)(\Delta_\mathcal{O})(2\Delta_\mathcal{O} - 1)}$$

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Can't resum...

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Focus on 3d so rotations are $SO(3)$

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Focus on 3d so rotations are $SO(3)$

Need to add potential for angular momentum:

$$F(\Delta_\phi, \Delta_\mathcal{O}; q, y) = q^{-\Delta_\mathcal{O}} \langle \mathcal{O} | \phi | \mathcal{O} \rangle^{-1} \text{Tr}_\mathcal{O}[\phi q^D y^{J_z}]$$

Care about $y = 1$.

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Care about $y = 1$.

Casimir operator of conformal group (in 3d)

$$C = D(D - 3) + J^2 - P_\mu K^\mu$$

has eigenvalue $\Delta(\Delta - 3)$ on irrep.

Differential equation

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Insert Casimir into trace

$$\Rightarrow \text{Tr}_{\mathcal{O}}[C\phi q^D y^{J_z}] \sim \Delta_{\mathcal{O}}(\Delta_{\mathcal{O}} - 3)F(\Delta_{\phi}, \Delta_{\mathcal{O}}; q, y)$$

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Convert to differential equation using algebra and

$$D \sim q\partial_q, \quad J_z \sim y\partial_y$$

$$\Rightarrow \mathcal{C}(q, y, x)F(\Delta_{\phi}, \Delta_{\mathcal{O}}; q, y) = \Delta_{\mathcal{O}}(\Delta_{\mathcal{O}} - 3)F(\Delta_{\phi}, \Delta_{\mathcal{O}}; q, y)$$

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Can't solve exactly...

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Euclidean AdS space is like
a cylinder

Anti de Sitter space

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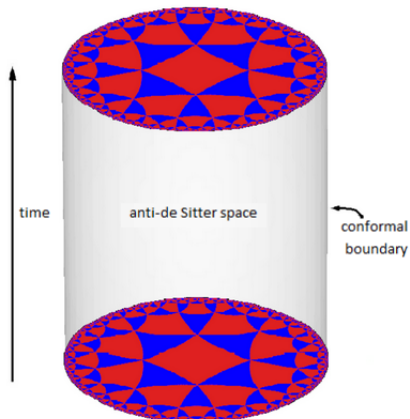
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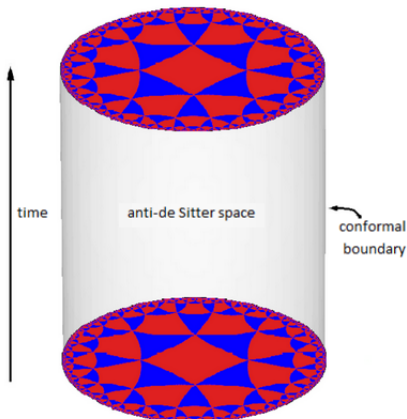
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Euclidean AdS space is like
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Isometry group of AdS_{d+1}
= Conformal group on \mathbb{R}^d



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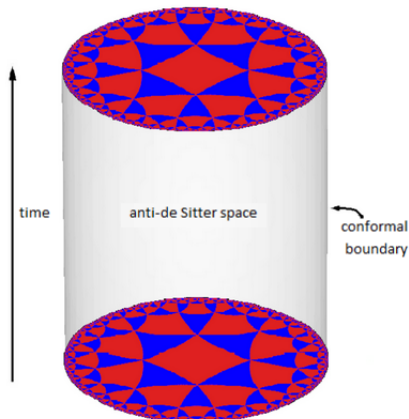
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Isometry group of AdS_{d+1}
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\Rightarrow Consider CFT on
boundary and compute
using bulk



Witten diagram

Thermal state so periodic in time ($t \sim t + \beta$)

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Witten diagram

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Interested in 1-point function of ϕ with one loop of \mathcal{O} (from $\langle \mathcal{O} | \phi | \mathcal{O} \rangle$)

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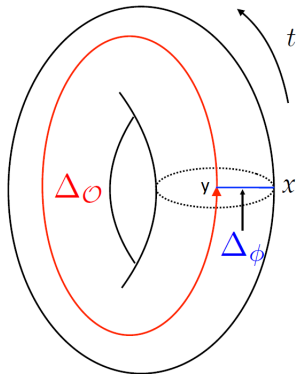
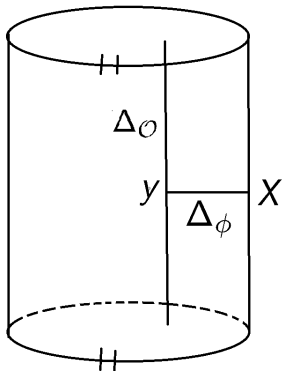
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Final result

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Explicitely (without angular potential)

$$F(\Delta_\phi, \Delta_\mathcal{O}; q) = \int_{AdS} dy \sqrt{g} G_{bulk-bdy}^{\Delta_\phi}(y, x) G_{bulk-bulk}^{\Delta_\mathcal{O}}(y, y_{thermal})$$

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$$F(\Delta_\phi, \Delta_{\mathcal{O}}; q) = \int_{AdS} dy \sqrt{g} G_{bulk-bdy}^{\Delta_\phi}(y, x) G_{bulk-bulk}^{\Delta_{\mathcal{O}}}(y, y_{thermal})$$

Can actually do the integral

$$\Rightarrow F(\Delta_\phi, \Delta_{\mathcal{O}}; q) = \frac{1}{(1-q)^{2\Delta_{\mathcal{O}}}} \times {}_3F_2 \left[\frac{1-d}{2} + \Delta_{\mathcal{O}}, \Delta_{\mathcal{O}} - \frac{\Delta_\phi}{2}, \frac{\Delta_\phi - d}{2} + \Delta_{\mathcal{O}}; \Delta_{\mathcal{O}}, -d + 2\Delta_{\mathcal{O}} + 1; -\frac{4q}{(q-1)^2} \right]$$

Conclusion and future work

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Next steps:

- Spinning operators
- Higher point functions
- Do bootstrap

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Thank you!